

# Earthquake Response of Inelastic Systems

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Motivation

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Cyclic Behavior of Structural Members

Cyclic Behavior

Elastic-plastic Idealization

E-P Idealization

Earthquake Response of E-P Systems

Earthquake  
Response of E-P  
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Normalized Equation of Motion

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Yielding

Effects of Yielding

Inelastic Response, different values of  $\bar{f}_y$

In Earthquake Engineering it is common practice to design against a large earthquake, that has a given mean period of return (say 500 years), quite larger than the expected life of the construction.

A period of return of 500 years means that in a much larger interval, say 50000 years, you expect say 100 earthquakes that are no smaller (in the sense of some metrics, e.g., the peak ground acceleration) than the design earthquake.

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For the unlikely occurrence of a large earthquake, a large damage in the construction can be deemed acceptable as far as

- ▶ no human lives are taken in a complete structural collapse and
- ▶ in the mean, the costs for repairing a damaged building are not disproportionate to its value.



# What to do?

To ascertain the amount of acceptable reduction of earthquake loads it is necessary to study

- ▶ the behavior of structural members and systems subjected to cyclic loading outside the elastic range, to understand the amount of plastic deformation and accumulated plastic deformation that can be sustained before collapse and
- ▶ the global structural behavior for inelastic response, so that we can relate the reduction in design parameters to the increase in members' plastic deformation.

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- ▶ the global structural behavior for inelastic response, so that we can relate the reduction in design parameters to the increase in members' plastic deformation.

The first part of this agenda pertains to Earthquake Engineering proper, the second part is across EE and Dynamics of Structures, and today's subject.

# Cyclic behavior

Investigation of the cyclic behavior of structural members, sub-assemblages and scaled or real sized building model, either in labs or via numerical simulations, constitutes the bulk of EE.

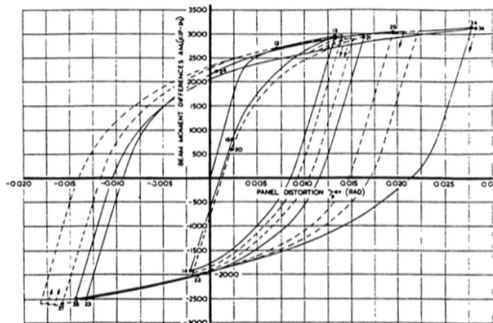
What is important, at the moment, is the understanding of how different these behaviors can be, due to different materials or structural configurations, with instability playing an important role.

We will see 3 different diagrams, force vs deformation, for a clamped steel beam subjected to flexion, a reinforced concrete sub-assemblage and a masonry wall.

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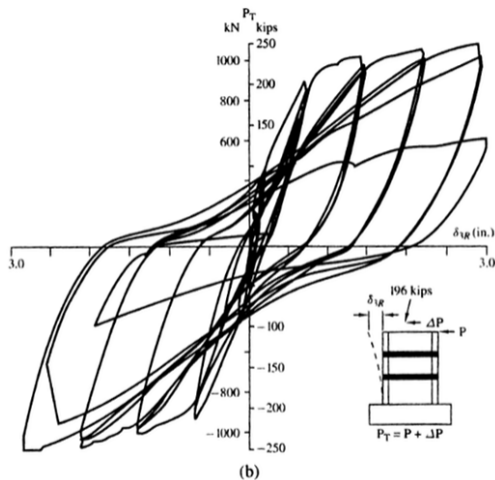
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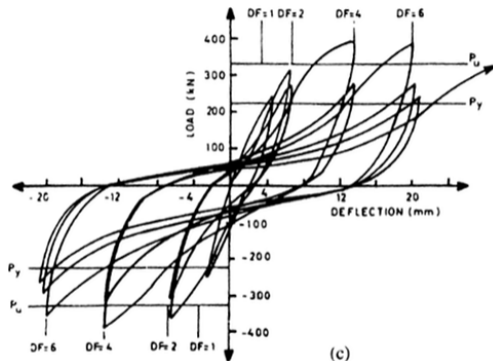
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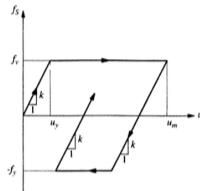
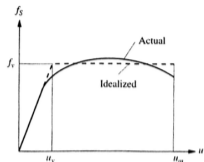
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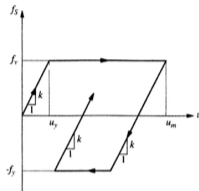
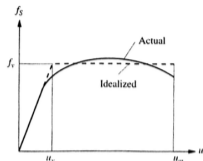
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# E-P model



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# E-P model

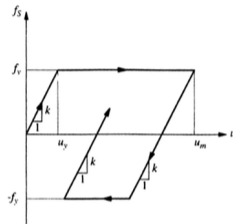
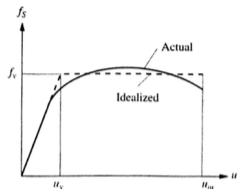


A more complex behavior may be represented with an elastic-perfectly plastic (e-p) bi-linear idealization, see figure, where two important requirements are obeyed

1. the initial stiffness of the idealized e-p system is the same of the real system, so that the natural frequencies of vibration for small deformation are equal,
2. the yielding strength is chosen so that the sum of stored and dissipated energy in the e-p system is the same as the energy stored and dissipated in the real system.



## E-P model, 2



In perfect plasticity, when yielding (a) the force is constant,  $f_S = f_y$  and (b) the stiffness is null,  $k_y = 0$ . The force  $f_y$  is the yielding force, the displacement  $x_y = f_y/k$  is the yield deformation.

In the right part of the figure, you can see that at unloading ( $dx = 0$ ) the stiffness is equal to the initial stiffness, and we have  $f_S = k(x - x_{p_{tot}})$  where  $x_{p_{tot}}$  is the total plastic deformation.

# Definitions

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$$\bar{f}_y = \min \left\{ \frac{f_y}{f_0} = \frac{x_y}{x_0}, 1 \right\}.$$

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**normalized spring force,  $\bar{f}_S$**  the ratio of the e-p spring force to the yield strength,

$$\bar{f}_S = f_S / f_y.$$

## Definitions, cont.

**equivalent acceleration,  $a_y$**  the (pseudo-)acceleration required to yield the system,  $a_y = \omega_n^2 x_y = f_y/m$ .

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*NB the ratio between the e-p and elastic peak responses is given by*

$$\frac{x_m}{x_0} = \frac{x_m}{x_y} \frac{x_y}{x_0} = \mu \bar{f}_y = \frac{\mu}{R_y} \rightarrow \mu = R_y \frac{x_m}{x_0}.$$

## Normalizing the force

We want to show that, for a given excitation  $\ddot{x}_g(t)$ , the response depends on 3 parameters,  $\omega_n = \sqrt{k/m}$ ,  $\zeta = c/(2\omega_n m)$  and  $x_y$ .

For an e-p system, the equation of motion (*EOM*) is

$$m\ddot{x} + c\dot{x} + f_S(x, \dot{x}) = -m\ddot{u}_g(t)$$

with  $f_S$  as shown in a previous slide. The *EOM* must be integrated numerically to determine the time history of the e-p response,  $x(t)$ . If we divide the *EOM* by  $m$ , recalling our definition of the normalized spring force, the last term is

$$\frac{f_S}{m} = \frac{1}{m} \frac{f_y}{f_y} f_S = \frac{1}{m} k x_y \frac{f_S}{f_y} = \omega_n^2 x_y \bar{f}_S$$

and we can write

$$\ddot{x} + 2\zeta\omega_n\dot{x} + \omega_n^2 x_y \bar{f}_S(x, \dot{x}) = -\ddot{u}_g(t)$$

# Normalizing the displacements

With the position  $x(t) = \mu(t) x_y$ , substituting in the *EOM* and dividing all terms by  $x_y$ , it is

$$\ddot{\mu} + 2\omega_n \zeta \dot{\mu} + \omega_n^2 \bar{f}_S(\mu, \dot{\mu}) = -\frac{\omega_n^2 \ddot{x}_g}{\omega_n^2 x_y} = -\omega_n^2 \frac{\ddot{x}_g}{a_y}$$

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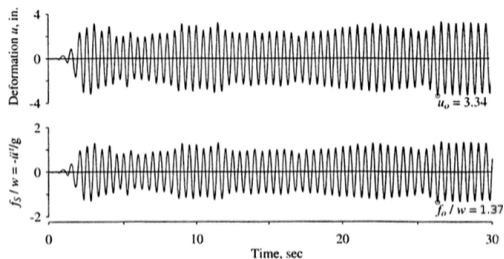
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The equivalent acceleration can be expressed in terms of the normalized yield strength  $\bar{f}_y$ ,

$$a_y = \frac{f_y}{m} = \frac{\bar{f}_y f_0}{m} = \frac{\bar{f}_y k x_0}{m} = \bar{f}_y \omega_n^2 x_0$$

and recognizing that  $x_0$  depends only on  $\zeta$  and  $\omega_n$  we conclude that, for given  $\ddot{x}_g(t)$  and  $\bar{f}_S(\mu, \dot{\mu})$  the ductility response depends only on  $\zeta$ ,  $\omega_n$ ,  $\bar{f}_y$ .

# Elastic response, required parameters

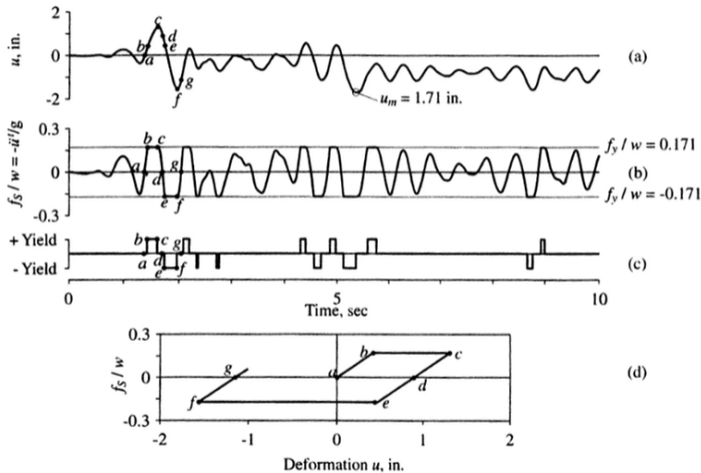


Response of linear system with  $T_n = 0.5$  sec and  $\zeta = 0$  to El Centro

In the figure above, the elastic response of an undamped,  $T_n = 0.5$  s system to the NS component of the El Centro 1940 ground motion (all our examples will be based on this input motion).

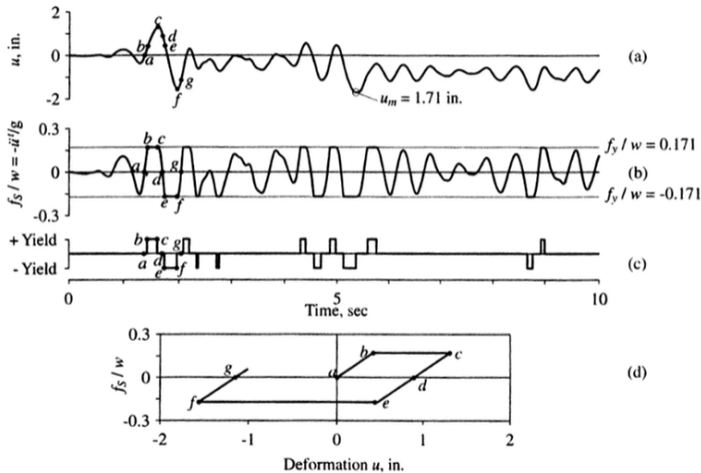
Top, the deformations, bottom the elastic force normalized with respect to weight, from the latter peak value we know that all e-p systems with  $f_y < 1.37w$  will experience plastic deformations during the EC1940NS ground motion.

# Inelastic response, $\bar{f}_y = 1/8$



The various response graphs above were computed for  $\bar{f}_y = 0.125$  (i.e.,  $R_y = 8$  and  $f_y = \frac{1.37}{8}w = 0.171w$ ) and  $\zeta = 0$ ,  $T_n = 0.5$  s.

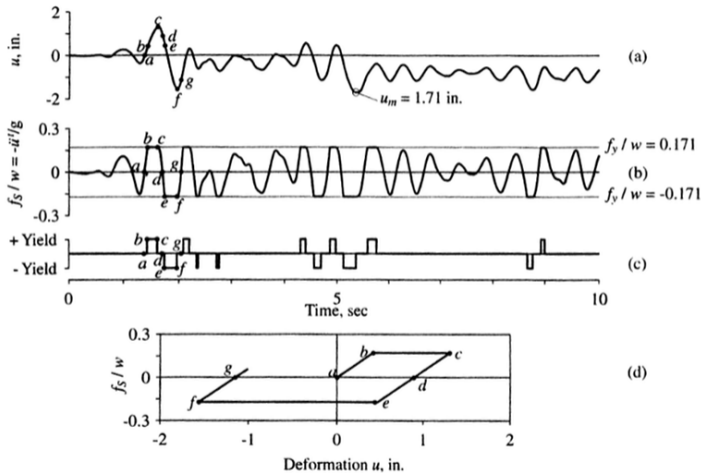
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Top, the deformation response, note that the peak response is  $x_m = 1.71$  in, different from  $x_0 = 3.34$  in; it is  $\mu = R_y \frac{x_m}{x_0} = 4.09$

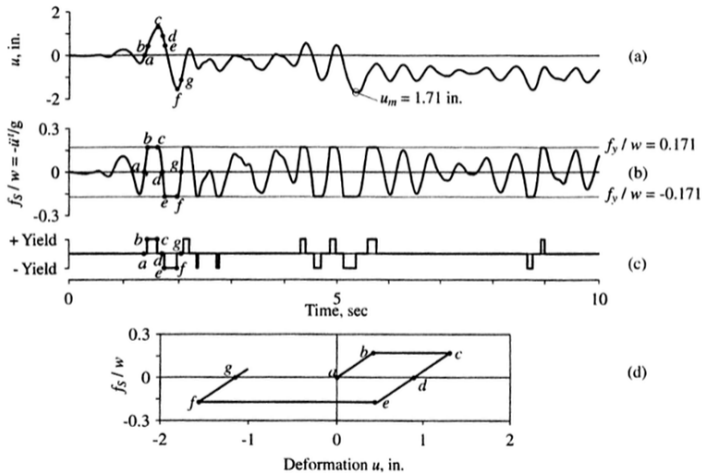


# Inelastic response, $\bar{f}_y = 1/8$



Second row, normalized force  $f_S/w$ , note that the response is *clipped* at  $f_S = f_y = 1.171w$

# Inelastic response, $\bar{f}_y = 1/8$



Third row, response in terms of yielding state, positive or negative depending on the sign of velocity

## Inelastic response, $\bar{f}_y = 1/8$

The force-deformation diagram for the first two excursions in plastic domain, the time points  $a, b, c, d, e, f$  and  $g$  are the same in all 4 graphs:

- ▶ until  $t = b$  we have an elastic behavior,
- ▶ until  $t = c$  the velocity is positive and the system accumulates positive plastic deformations,
- ▶ until  $t = e$  we have an elastic unloading (note that for  $t = d$  the force is zero, the deformation is equal to the total plastic deformation),
- ▶ until  $t = f$  we have another plastic excursion, accumulating negative plastic deformations
- ▶ until at  $t = f$  we have an inversion of the velocity and an elastic reloading.

## Response for different $\bar{f}_y$ 's

$\bar{f}_y$	$x_m$	$x_{perm}$	$\mu$
1.000	2.25	0.00	1.00
0.500	1.62	0.17	1.44
0.250	1.75	1.10	3.11
0.125	2.07	1.13	7.36

This table was computed for  $T_n = 0.5$  s and  $\zeta = 5\%$  for the EC1940NS excitation.

Elastic response was computed first, with peak response  $x_0 = 2.25$  in and peak force  $f_0 = 0.919w$ , later the computation was repeated for  $\bar{f}_y = 0.5, 0.25, 0.125$ .

In our example, all the peak values of the e-p responses are smaller than the elastic one, but this is not always true, and shouldn't be generalized. The permanent displacements increase for decreasing yield strengths, and also this fact shouldn't be generalized.

Last, the ductility ratios increase for decreasing yield strengths, for our example it is  $\mu \approx R_y$ .

# Ductility demand and capacity

We can say that, for a given value of the normalized yield strength  $\bar{f}_y$  or of the yield strength reduction factor  $R_y$ , there is a *ductility demand*, a measure of the extension of the plastic behavior that is required when we reduce the strength of the construction.

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Corresponding to this ductility demand our structure must be designed so that there is a sufficient *ductility capacity*.

Ductility capacity is, in the first instance, the ability of individual members to sustain the plastic deformation demand without collapsing, the designer must verify that the capacity is greater than the demand for all structural members that go non linear during the seismic excitation.

Motivation

Cyclic Behavior

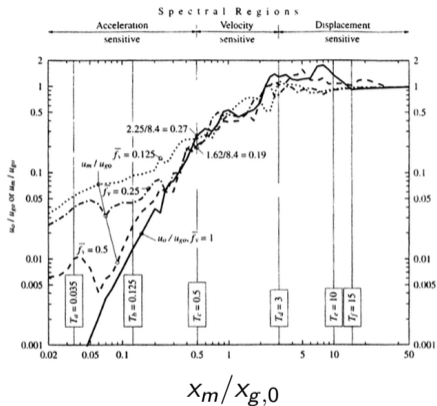
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# Effects of $T_n$

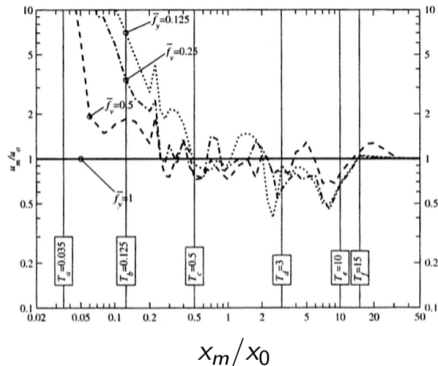


For EC1940NS, for  $\zeta = .05$ , for different values of  $T_n$  and for  $\bar{f}_y = 1.0, 0.5, 0.25, 0.125$  the peak response  $x_0$  of the equivalent system (in black) and the peak responses of the 3 inelastic systems has been computed.

There are two distinct zones: left there is a strong dependency on  $\bar{f}_y$ , the peak responses grow with  $R_y$ ; right the 4 curves intersect one with the others and there is no clear dependency on  $\bar{f}_y$ .



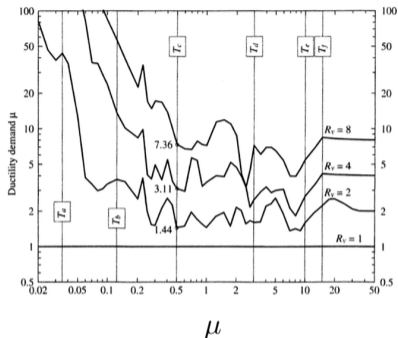
# Effects of $T_n$



With the same setup as before, here it is the ratio of the  $x_m$ 's to  $x_0$ , what is evident is the fact that, for large  $T_n$ , this ratio is equal to 1... this is justified because, for large  $T_n$ 's, the mass is essentially at rest, and the deformation, either elastic or elastic-plastic, are equal and opposite to the ground displacement.

Also in the central part, where elastic spectrum ordinates are dominated by the ground velocity, there is a definite tendency for the  $x_m/x_0$  ratio, that is  $x_m/x_0 \approx 1$

# Effects of $T_n$

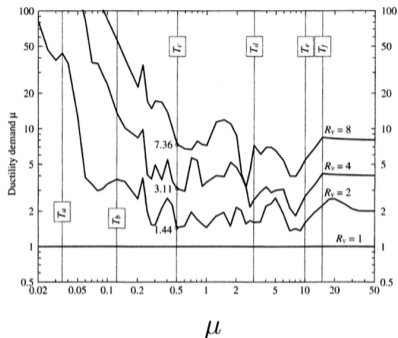


With the same setup as before, in this graph are reported the values of the ductility factor  $\mu$ .

The values of  $\mu$  are almost equal to  $R_y$  for large values of  $T_n$ , and in the limit, for  $T_n \rightarrow \infty$ , there is a strict equality. An even more interesting observation regard the interval  $T_c \leq T_n \leq T_f$ , where the values of  $\mu$  oscillate near the value of  $R_y$ .

On the other hand, the behavior is completely different in the acceleration controlled zone, where  $\mu$  grows very fast, and the ductility demand is very high even for low values (0.5) of the yield strength reduction factor.

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The results we have discussed are relative to one particular excitation, nevertheless research and experience confirmed that these propositions are true also for different earthquake records, taking into account the differences in the definition of spectral regions.

The first step in an anti seismic design is to set an available ductility (based on materials, conception, details).

In consequence, we desire to know the yield displacement  $u_y$  or the yield force  $f_y$

$$f_y = ku_y = m\omega_n^2 u_y$$

for which the ductility demand imposed by the ground motion is not greater than the available ductility.

# Response Spectrum for Yield States

For each  $T_n$ ,  $\zeta$  and  $\mu$ , the *Yield-Deformation Response Spectrum* ( $D_y$ ) ordinate is the corresponding value of  $u_y$ :  $D_y = u_y$ . Following the ideas used for Response and Design Spectra, we define  $V_y = \omega_n u_y$  and  $A_y = \omega_n^2 u_y$ , that we will simply call pseudo-velocity and pseudo-acceleration spectra.

Using our definitions of pseudo acceleration, we can find a more significant expression for the design force:

$$f_y = k u_y = m \omega_n^2 u_y = m A_y = w \frac{A_y}{g},$$

where  $w$  is the weight of the structure.

# Response Spectrum for Yield States

For each  $T_n$ ,  $\zeta$  and  $\mu$ , the *Yield-Deformation Response Spectrum* ( $D_y$ ) ordinate is the corresponding value of  $u_y$ :  $D_y = u_y$ . Following the ideas used for Response and Design Spectra, we define  $V_y = \omega_n u_y$  and  $A_y = \omega_n^2 u_y$ , that we will simply call pseudo-velocity and pseudo-acceleration spectra.

Using our definitions of pseudo acceleration, we can find a more significant expression for the design force:

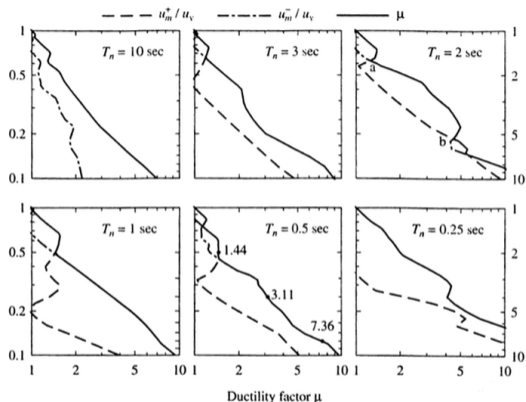
$$f_y = k u_y = m \omega_n^2 u_y = m A_y = w \frac{A_y}{g},$$

where  $w$  is the weight of the structure.

Our definition of inelastic spectra is compatible with the definition of elastic spectra, because for  $\mu = 1$  it is  $u_y = u_0$ .

Finally, the  $D_y$  spectrum and its derived pseudo spectra can be plotted on the tripartite log-log graph.

# Computing $D_y$



On the left, for different  $T_n$ 's and  $\mu = 5\%$ . the independent variable is in the ordinates, either  $\bar{f}_y$  (left) or  $R_y$  (right) the strength reduction factor. Dash-dot lines is  $u_m^+ / u_y$ , dash-dot is  $u_m^- / u_y$ .  $u_m^+$  and  $u_m^-$  are the peaks of positive and negative displacements of the inelastic system. The maximum of their ratios to  $u_y$  is the ductility  $\mu$ .

If we look at these graphs using  $\mu$  as the independent variable, it is possible that for a single value of  $\mu$  there are different values on the tick line: in this case, for security reasons, the designer must design for the higher value of  $\bar{f}_y$ .

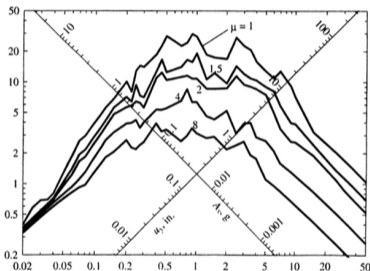
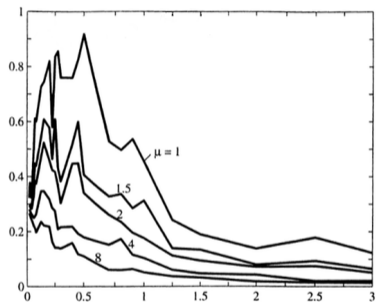
# Example

For EC1940NS,  $z = 5\%$ , the yield-strength response spectra for  $\mu = 1.0, 1.5, 2.0, 4.0, 8.0$ .



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On the left, a lin-lin plot of the pseudo-acceleration normalized (and dimensionless) with respect to  $g$ , the acceleration of gravity.

On the right, a log-log tripartite plot of the same spectrum.

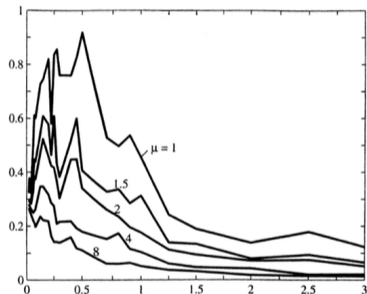
Even a small value of  $\mu$  produces a significant reduction in the required strength.

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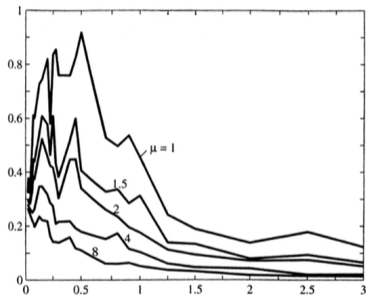
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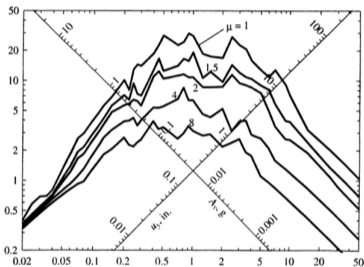
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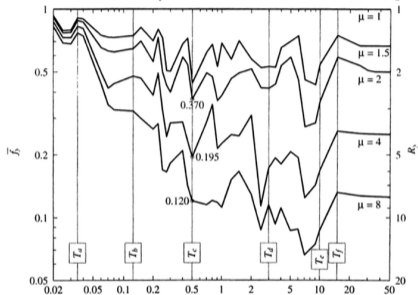
A lin-lin plot of the pseudo-acceleration normalized (and dimensionless) with respect to  $g$ , the acceleration of gravity.



A log-log tripartite plot of the same spectrum. Even a small value of  $\mu$  produces a significant reduction in the required strength.

# $\bar{f}_y$ vs $\mu$

We have seen that  $\bar{f}_y = \bar{f}_y(\mu, T_n, \zeta)$  is a monotonically increasing function of  $\mu$  for fixed  $T_n$  and  $\zeta$ .

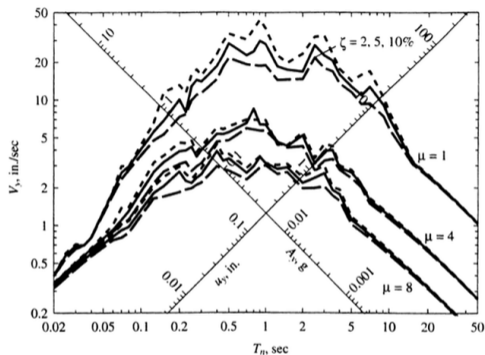


Left the same spectra of the previous slides, plotted in a different format,  $\bar{f}_y$  vs  $T_n$  for different values of  $\mu$ .

The implication of this figure is that an anti seismic design can be based on strength, ductility or a combination of the two.

For  $T_n = 1.0$ , the peak force for EC1940NS in an elastic system is  $f_0 = 0.919w$ , so it is possible to design for  $\mu = 1.0$ , hence  $f_y = 0.919w$  or for an high value of ductility,  $\mu = 8.0$ , hence  $f_y = 0.120 \cdot 0.919w$  or. If  $\mu = 8.0$  is hard to obtain, one can design for  $\mu = 4.0$  and a yielding force of 0.195 times  $f_0$ .

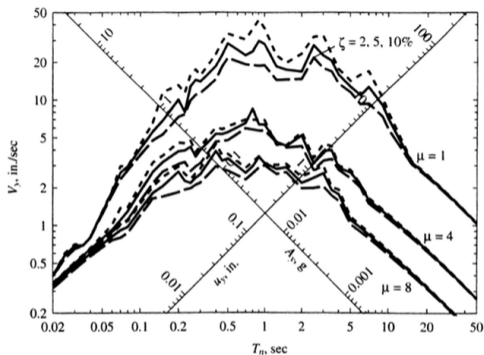
# Yielding and Damping



El Centro 1940 NS, elastic response spectra and inelastic spectra for  $\mu = 4$  and  $\mu = 8$ , for different values of  $\zeta$  (2%, 5% and 10%).

Effects of damping are relatively important and only in the velocity controlled area of the spectra, while effects of ductility are always important except in the high frequency range.

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Overall, the ordinates reduction due to modest increases in ductility are

$$\int^{x(t)} m \ddot{x} dx + \int^{x(t)} c \dot{x} dx + \int^{x(t)} f_S(x, \dot{x}) dx = - \int^{x(t)} m \ddot{x}_g dx$$

This is an energy balance, between the input energy  $\int m \ddot{x}_g$  and the sum of the kinetic, damped, elastic and dissipated by yielding energy.

In every moment, the elastic energy  $E_S(t) = \frac{f_S^2(t)}{2k}$  so the yielded energy is

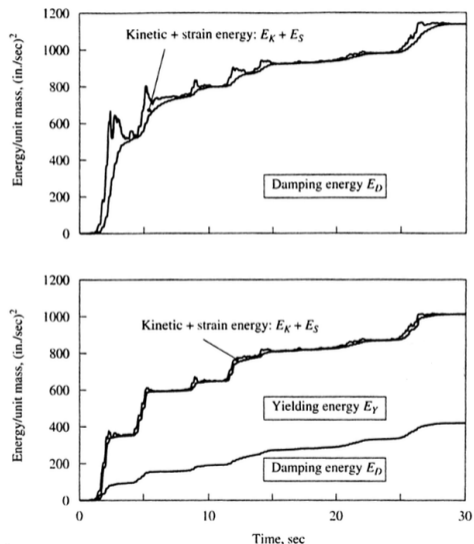
$$E_Y = \int f_S(x, \dot{x}) dx - \frac{f_S^2(t)}{2k}.$$

The damped energy can be written as a function of  $t$ , as  $dx = \dot{x} dt$ :

$$E_D = \int c \dot{x}^2(t) dt$$



# Energy Dissipation



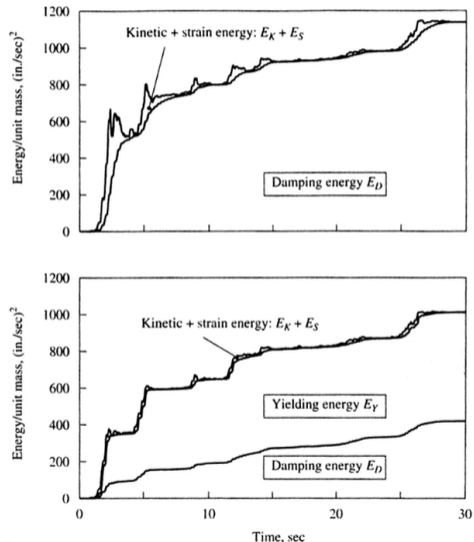
For a system with  $m = 1$  and

a)  $\bar{f}_y = 1$

b)  $\bar{f}_y = 0.25$

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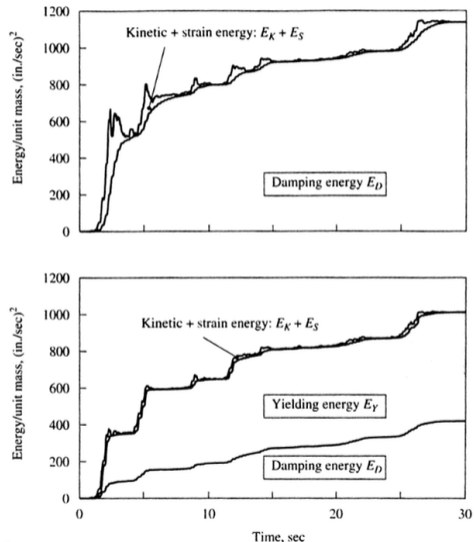
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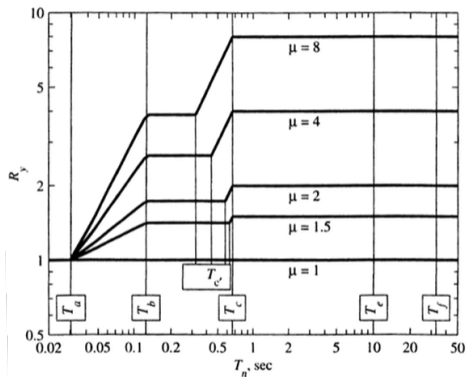
In b), yielding energy is dissipated by means of some structural damage.

Two possible approaches:

1. compute response spectra for constant ductility demand for many consistent records, compute response parameters statistics and derive inelastic design spectra from these statistics, as in the elastic design spectra procedures;
2. directly modify the elastic design spectra to account for the ductility demand values.

The first procedure is similar to what we have previously seen, so we will concentrate on the second procedure, that it is much more used in practice.

# $R_y - \mu - T_n$ equations

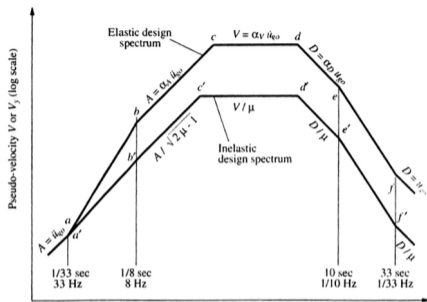


Based on observations and energetic considerations, the plots of  $R_y$  vs  $T_n$  for different  $\mu$  values can be approximated with straight lines in a log-log diagram, where the constant pieces are defined in terms of the key periods in  $D - V - A$  graphs.

$$R_y = \begin{cases} 1 & T_n < T_a \\ \sqrt{2\mu - 1} & T_b < T_n < T_c' \\ \mu & T_c < T_n \end{cases}$$

The key period  $T_{c'}$  is different from  $T_c$ , as we will see in the next slide; the constant pieces are joined with straight lines in the log-log diagram.

# Construction of Design Spectrum



Start from a given elastic design spectrum, defined by the points a-b-c-d-e-f.

Choose a value  $\mu$  for the ductility demand.

Reduce all ordinates right of  $T_c$  by the factor  $\mu$ , reduce the ordinates in the interval  $T_b < T_n < T_c$  by  $\sqrt{2\mu/1}$ .

Draw the two lines  $A = \frac{\alpha_A \ddot{x}_{g0}}{\sqrt{2\mu - 1}}$  and  $V = \frac{\alpha_V \dot{x}_{g0}}{\mu}$ , their intersection define the key point  $T_{c'}$ .

Connect the point  $(T_a, A = \ddot{x}_{g0})$  and the point  $(T_b, A = \frac{\alpha_V \dot{x}_{g0}}{\mu})$  with a straight line.

As we already know (at least in principles) the procedure to compute the elastic design spectra for a given site from the peak values of the ground motion, using this simple procedure it is possible to derive the inelastic design spectra for any ductility demand level.

# Important Relationships

For different zones on the  $T_n$  axis, the simple relationships we have previously defined can be made explicit using the equations that define  $R_y$ , in particular we want relate  $u_m$  to  $u_0$  and  $f_y$  to  $f_0$  for the elastic-plastic system and the equivalent system.

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Similar equations can be established also for the inclined connection segments in the  $R_y$  vs  $T_n$  diagram.

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- ▶ The design peak deformation,  $u_m = \mu D_y / R_y$ , is

$$u_m = \frac{\mu}{R_y(\mu, T_n)} \frac{A_y}{\omega_n^2}.$$

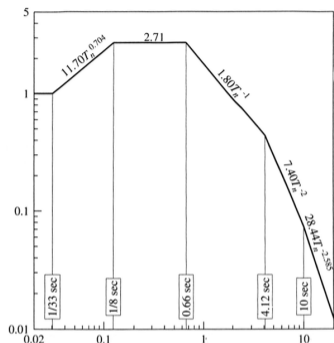


# Example

One storey frame, weight  $w$ , period is  $T_n = 0.25$  s, damping ratio is  $\zeta = 5\%$ , peak ground acceleration is  $\ddot{x}_{g0} = 0.5$  g. Find design forces for

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# Example



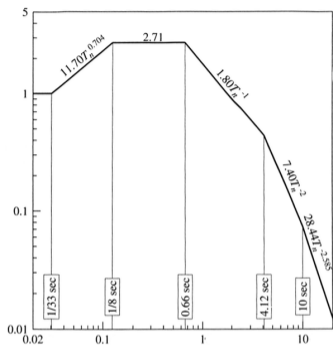
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In the figure, a reference elastic spectrum for  $\ddot{x}_{g0} = 1$  g,  $A_y(0.25) = 2.71$  g; for  $\ddot{x}_{g0} = 0.5$  g it is  $f_0 = 1.355w$ .

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For  $T_n = 0.25$  s,  $R_y = \sqrt{2\mu - 1}$ , hence

$$f_y = \frac{1.355w}{\sqrt{2\mu - 1}}, \quad u_m = \frac{\mu}{\sqrt{2\mu - 1}} \frac{A_y}{\omega_n^2} = \frac{\mu}{\sqrt{2\mu - 1}} \frac{1.355g T_n^2}{4\pi^2}.$$

$$\mu = 1 : \quad f_y = 1.355w, \quad u_m = 2.104 \text{ cm};$$

$$\mu = 4 : \quad f_y = 0.512w, \quad u_m = 3.182 \text{ cm};$$

$$\mu = 8 : \quad f_y = 0.350w, \quad u_m = 4.347 \text{ cm}.$$