

## 2 DOF System

```
from numpy import array, cos, diag, exp, linspace, pi, poly1d, sin, sqrt
from scipy.linalg import inv, det, eigh
import matplotlib.pyplot as plt
from IPython.display import Latex
t = linspace(0, 20, 501)
```

## First Part, Structural Matrices & Eigenvalues

Let's start with a few helper functions

```
def p(*l): return poly1d(l)
def i(p1, p2):
    p = (p1*p2).integ()
    return(p(1)-p(0))
def pmat(mat, fmt='%.3f'):
    inside = r'\'.join('&'.join(fmt%x for x in row) for row in mat)
    return r'\begin{bmatrix}' + inside + r'\end{bmatrix}'
def dl(s): display(Latex('$$'+s+'$$'))
```

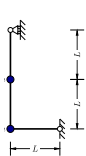


Figure 1a.

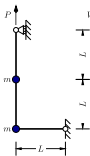


Figure 1b.

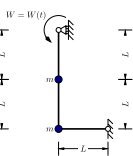


Figure 1c.

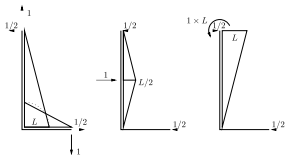
The beam is split in three intervals of definition for the bending moments, all intervals of length  $L$ : **1** from the top to the middle of the vertical part, **2** from the corner to the middle of the vertical part, **3** from the hinge to the corner.

The bending moments, for a 1 vertical load, for a 1 horizontal load at midspan and for a  $1 \times L$  couple applied to the upper support are

vertical:  $+\frac{1}{2}x + 0, \quad -\frac{1}{2}x + L, \quad +1x + 0;$

horizontal:  $+\frac{1}{2}x + 0, \quad +\frac{1}{2}x + 0, \quad +0x + 0;$

couple  $-\frac{1}{2}x + L, \quad +\frac{1}{2}x + 0, \quad +0x + 0.$



The bending moments are here expressed as polynomials ( $p(\dots)$ ) and collected in a list for each load conditions, those lists collected in an outer list  $m$ .

```
m = [  
    [p( 0.5, 0), p(-0.5, 1), p( 1, 0)],  
    [p( 0.5, 0), p( 0.5, 0), p( 0, 0)],  
    [p(-0.5, 1), p( 0.5, 0), p( 0, 0)],  
]
```

The bending moments due to unit forces are used to compute the  $3 \times 3$  flexibility matrix,  $c$  stands for curvatures and  $bm$  for bending moments.

```
F33 = array([[sum(i(ci, bmi) for ci, bmi in zip(c, bm)) for c in m] for bm in m])
```

Also needed are the  $2 \times 2$  flexibility, the  $3 \times 3$  stiffness matrix and the  $2 \times 2$  one (note that, having the flexibility, we used a procedure different from the static condensation); the  $3 \times 3$  matrix is needed to compute the equivalent load on the dynamic degrees of freedom below.

```
F22 = F33[ :-1, :-1]
K33 = inv(F33)
K22 = inv(F22)
```



```

d1(r'F33=\frac{L^3}{12EJ}' + pmat(F33*12) + ', \quad F22= \frac{L^3}{12EJ}' + pmat(F22*12))
d1('K33=\frac{3EJ}{13L^3}' + pmat(K33*13/3, fmt='% .3f') + ', \quad K22=\frac{4EJ}{5L^3}' + pmat(K22*5/4, fmt='% .3f'))

```

$$F_{33} = \frac{L^3}{12EJ} \begin{bmatrix} 12.000 & 3.000 & 4.000 \\ 3.000 & 2.000 & 3.000 \\ 4.000 & 3.000 & 8.000 \end{bmatrix},$$

$$F_{22} = \frac{L^3}{12EJ} \begin{bmatrix} 12.000 & 3.000 \\ 3.000 & 2.000 \end{bmatrix}$$

$$K_{33} = \frac{3EJ}{13L^3} \begin{bmatrix} 7.000 & -12.000 & 1.000 \\ -12.000 & 80.000 & -24.000 \\ 1.000 & -24.000 & 15.000 \end{bmatrix},$$

$$K_{22} = \frac{4EJ}{5L^3} \begin{bmatrix} 2.000 & -3.000 \\ -3.000 & 12.000 \end{bmatrix}$$

Eventually, the mass matrix (note that a vertical motion involves *both* the supported masses).

```
M22 = array(((2.0,0.0),(0.0,1.0)))  
d1('M22 = m' + pmat(M22, fmt='%1f'))
```

$$M_{22} = m \begin{bmatrix} 2.0 & 0.0 \\ 0.0 & 1.0 \end{bmatrix}$$

Let's compute the eigenvalues `evals`, the eigenvectors `evecs` (aliased to `Psi`) and the circular frequencies. While at it, we *need* the inverse of the starred mass matrix, don't we?

```
evals, evecs = eigh(K22, M22)
Psi = evecs
cfreqs = sqrt(evals)
Mstarinv = array(((1.0,0),(0,1.0)))
dl(r'\boldsymbol{\Lambda}^2=\omega_0^2'+pmat(diag(evals)))
dl(r'\boldsymbol{\Lambda} = \omega_0'+pmat(diag(sqrt(evals))))
dl('\boldsymbol{\Psi}'+pmat(evecs))
```

$$\mathbf{\Lambda}^2 = \omega_0^2 \begin{bmatrix} 0.484 & 0.000 \\ 0.000 & 9.916 \end{bmatrix}$$

$$\mathbf{\Lambda} = \omega_0 \begin{bmatrix} 0.696 & 0.000 \\ 0.000 & 3.149 \end{bmatrix}$$

$$\mathbf{\Psi} = \begin{bmatrix} -0.695 & -0.129 \\ -0.183 & 0.983 \end{bmatrix}$$

## Second Part, Free Vibrations

The system, after the load removal, is freely vibrating so that the modal responses are simply cosines, whose amplitude is given by the components of the  $q_0$  vector.

The initial displacements can be computed from the static load vector, next we convert to modal coordinates and check our result

```
x0 = F22@array((1, 0)) # L**3/EJ * F * {1,0} EJ/300/L**2 = L/300 * F *{1,0}
d1('\frac{300\\,x_0}L = '+pmat(x0[:,None]))
q0 = Mstarinv@evecs.T@M22@x0
d1('\frac{300\\,q_0}L='+pmat(q0[:,None]))
d1(r'\Psi\, q_0 = \frac{L}{300}\,'+pmat((evecs@q0)[:,None], fmt="%g"))
```

$$\frac{300 x_0}{L} = \begin{bmatrix} 1.000 \\ 0.250 \end{bmatrix}$$

$$\frac{300 q_0}{L} = \begin{bmatrix} -1.436 \\ -0.013 \end{bmatrix}$$

$$\Psi q_0 = \frac{L}{300} \begin{bmatrix} 1 \\ 0.25 \end{bmatrix}$$

## Analytical expressions of modal responses

```
d1(r'\frac{300\,q_1(t)}{L} = %f \cos %.3f\, \omega_0 t' %(q0[0],cfreqs[0]))  
d1(r'\frac{300\,q_2(t)}{L} = %f \cos %.3f\, \omega_0 t' %(q0[1],cfreqs[1]))
```

$$\frac{300 q_1(t)}{L} = -1.436081 \cos 0.696 \omega_0 t$$

$$\frac{300 q_2(t)}{L} = -0.013051 \cos 3.149 \omega_0 t$$

## Analytical expressions of displacements

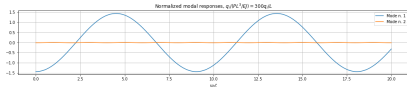
```
d1(r'\frac{300\,x_1(t)}{L} = %f \cos %.3f\,\omega_0 t %f \cos%.3f\,\omega_0 t'  
%(Psi[0,0]*q0[0], cfreqs[0], Psi[0,1]*q0[1], cfreqs[1]))  
d1(r'\frac{300\,x_2(t)}{L} = %f \cos %.3f\,\omega_0 t %f \cos%.3f\,\omega_0 t'  
%(Psi[1,0]*q0[0], cfreqs[0], Psi[1,1]*q0[1], cfreqs[1]))
```

$$\frac{300 x_1(t)}{L} = +0.998311 \cos 0.696 \omega_0 t + 0.001689 \cos 3.149 \omega_0 t$$

$$\frac{300 x_2(t)}{L} = +0.262831 \cos 0.696 \omega_0 t - 0.012831 \cos 3.149 \omega_0 t$$

## Plot of normalized modal responses

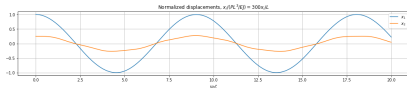
```
plt.rcParams["figure.figsize"] = (18,3)
lines = plt.plot(t, q0*cos(t[:,None]*cfreqs))
for n, line in enumerate(lines, 1): line.set_label('Mode n. %d'%n)
plt.legend(loc='best') ; plt.xlabel('\omega_0t$') ; plt.grid()
plt.title('Normalized modal responses, $q_i/(PL^3/EJ)=300q_i/L$');
```





## Plot of normalized displacements

```
lines = plt.plot(t, (q0*cos(t[:,None]*cfreqs))@evecs.T)
for n, line in enumerate(lines): line.set_label('$x_{%d}$'%(n+1))
plt.legend(loc='best') ; plt.xlabel('$\omega t$') ; plt.grid()
plt.title('Normalized displacements, $x_i/(PL^3/EJ)=300x_i/L$');
```



## **Third Part, Forced Response**

## The equivalent load

```
Kaa, Kab = K33[:2,:2], K33[:2,2:]
Kba, Kbb = K33[2:,:2], K33[2:,2:]

W = 1.0 # EJ/(400L)
p = -Kab@inv(Kbb)*W # Kab multiplies by L, inv(Kbb) divides by L**2,
                    # the result is the prod of adimensional matrix * W/L
p = p.flatten()    # p.flatten remedies an implementation detail
dl(r'p = \frac{EJ}{400\,L^2}' + pmat(p[:,None]))
pstar = evecs.T@p
dl(r'p^star = \frac{EJ}{400\,L^2}' + pmat(pstar[:,None]))
```

$$p = \frac{EJ}{400 L^2} \begin{bmatrix} -0.067 \\ 1.600 \end{bmatrix}$$

$$p^* = \frac{EJ}{400 L^2} \begin{bmatrix} -0.246 \\ 1.582 \end{bmatrix}$$

The frequency of the excitation is  $\omega_0$ , we write  $p_i^* = \pi_i \frac{EJ}{400 L^2}$

$$m\ddot{q}_i + m\omega_i^2 q_i = \frac{EJ}{400 L^2} \pi_i \sin \omega_0 t = \frac{EJ}{400 L^2} \frac{mL^3}{EJ} \omega_0^2 \pi_i \sin \omega_0 t$$

Dividing by  $m$  and simplifying, the EOM can be written as

$$\ddot{q}_i + \omega_i^2 q_i = \frac{L}{400} \omega_0^2 \pi_i \sin \omega_0 t = \delta \omega_0^2 \pi_i \sin \omega_0 t$$

The particular integral is  $\xi_i = C_i \sin \omega_0 t$ , substituting in the EOM

$$C_i(\omega_i^2 - \omega_0^2) = \delta \omega_0^2 \pi_i \quad \rightarrow \quad C_i = \frac{\omega_0^2}{\omega_i^2 - \omega_0^2} \delta \pi_i = \frac{\pi_i}{\lambda_i^2 - 1} \delta$$

`C = pstar/(evals-1)`

## Analytical expressions of modal responses

$$q_i = \delta C_i (\sin \omega_0 t - \beta_i \sin \omega_i t)$$

```
d1(r'\frac{400\,q_1}{L}=f(\sin\omega_0t+f\sin%.3f\,\omega_0t)'%(C[0],-1/cfreqs[0],cfreqs[0]))  
d1(r'\frac{400\,q_2}{L}=f(\sin\omega_0t+f\sin%.3f\,\omega_0t)'%(C[1],-1/cfreqs[1],cfreqs[1]))
```

$$\frac{400 q_1}{L} = +0.477752(\sin \omega_0 t - 1.437296 \sin 0.696 \omega_0 t)$$

$$\frac{400 q_2}{L} = +0.177391(\sin \omega_0 t - 0.317565 \sin 3.149 \omega_0 t)$$

## Analytical expressions of displacements

$$x_i = \sum_j \psi_{ij} q_j$$

```
for i in (0, 1):
    d1(r'x_%d = %f\sin\omega_0t %f\sin%.3f\,\omega_0t %f\sin%.3f\,\omega_0t'%(
        i+1, Psi[i,0]*C[0]+Psi[i,1]*C[1],
        -Psi[i,0]*C[0]/cfreqs[0], cfreqs[0],
        -Psi[i,1]*C[1]/cfreqs[1], cfreqs[1]))
```

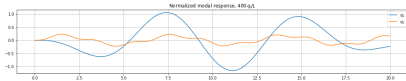
$$x_1 = -0.355072 \sin \omega_0 t + 0.477348 \sin 0.696 \omega_0 t \\ + 0.007290 \sin 3.149 \omega_0 t$$

$$x_2 = +0.086957 \sin \omega_0 t + 0.125674 \sin 0.696 \omega_0 t \\ - 0.055382 \sin 3.149 \omega_0 t$$

## Plot of normalized modal responses

```
q1 = C[0]*(sin(t)-sin(cfreqs[0]*t)/cfreqs[0])
q2 = C[1]*(sin(t)-sin(cfreqs[1]*t)/cfreqs[1])

plt.plot(t, q1, label='$q_1$')
plt.plot(t, q2, label='$q_2$')
plt.title(r'Normalized modal response, $400\backslash,q_i/L$')
plt.legend(loc='best') ; plt.grid() ;
```



## Plot of normalized displacements

```
x1 = Psi[0,0]*q1 + Psi[0,1]*q2
x2 = Psi[1,0]*q1 + Psi[1,1]*q2

plt.plot(t, x1, label='$x_1$')
plt.plot(t, x2, label='$x_2$')
plt.title(r'Normalized displacements, $400\backslash,x_i/L$')
plt.legend(loc='best') ; plt.grid();
```

